Supplemental Material for "Catch-Disperse-Release Readout for Superconducting Qubits"

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Evolution of Wigner function and probability distributions

In this supplemental material, we present the evolution of the Wigner function for resonator field corresponding to initial states $|00\rangle$ and $\overline{|10\rangle}$. Here the eigenstate $\overline{|10\rangle} = \cos \theta_1 |00\rangle - \sin \theta_1 |01\rangle$ with $\tan \theta_1 = 2g(\Delta_0^2 + 4g^2)^{-1/2}$, where Δ_0 is the initial qubit-resonator detuning and g is the qubit-resonator field coupling. Note that for large Δ_0 , the eigenstate $\overline{|10\rangle} \approx |10\rangle$ (in $|nm\rangle$, n represents qubit state and m represents the Fock state). We also present the evolution of the probability distributions of the measurement result x_{φ} at different times, and the waveform used for qubit frequency and for the microwave pulse. In the results presented up to page 14, we assumed a two-level model of the qubit.

We compute the Wigner function using the formula

$$W(\alpha) = Tr\left[D(-\alpha)\rho D(\alpha)e^{i\pi a^{\dagger}a}\right],$$

where $D(\alpha) = e^{-|\alpha|^2} e^{\alpha a^{\dagger}} e^{-\alpha^* a}$ is the usual displacement operator, ρ is the density operator for the resonator field, and a^{\dagger} and a are creation and annihilation operators for the resonator field.

The qubit frequency and the microwave pulse used in all the results shown in the following slides have the form

$$\omega_q(t) = \omega_0 + \frac{(\Delta_0 - \Delta)}{2} \left[\operatorname{Erf}\left(\frac{t - t_q}{\sqrt{2}\sigma_q}\right) - \operatorname{Erf}\left(\frac{t - t_{qe}}{\sqrt{2}\sigma_{qe}}\right) \right]$$
$$B(t) = \frac{B_0}{2} \left[\operatorname{Erf}\left(\frac{t - t_B}{\sqrt{2}\sigma_B}\right) - \operatorname{Erf}\left(\frac{t - t_B - \tau_B}{\sqrt{2}\sigma_B}\right) \right]$$

The parameters used for the plots presented in all the slides are:

$$\frac{\omega_0}{2\pi} = 6 \text{GHz}, \frac{\Delta_0}{2\pi} = 1 \text{GHz}, \frac{\Delta}{2\pi} = 80 \text{MHz}, \frac{g}{2\pi} = 30 \text{MHz}$$
$$\sigma_q = 4 \text{ns}, \sigma_{qe} = 1 \text{ns}, t_q = 10 \text{ns}, t_{qe} = 160 \text{ns}$$
$$\frac{B_0}{2\pi} = 0.4974 \text{GHz} (\sim 9 \text{ photons}), \sigma_B = 1 \text{ns}, t_B = 3 \text{ns}, \tau_B = 1 \text{ns}$$



 W_1 —Wigner function for resonator state corresponding to an initial state $|10\rangle$



 P_{0L} — probability for qubit to be in state $|0\rangle$ P_{0R} — probability for qubit to be in state $|1\rangle$

Initial state: $\overline{|10\rangle}$ (qubit in state $|1\rangle$)

 P_{1R} — probability for qubit to be in state $|1\rangle$ P_{1L} — probability for qubit to be in state $|0\rangle$

























Effect of level $|2\rangle$ on the evolution of the Wigner function of the resonator field

In the previous slides, we assumed the two-level model of the qubit. Real superconducting qubits, however, are only slightly anharmonic oscillators, thus the effect of the next excited level $|2\rangle$ is often important as shown the following slides.

The qubit anharmonicity is given by $\mathcal{A} = \omega_q - \omega_{q,12}$. For phase and transmon superconducting qubits the anharmonicity is about 3%. Here we assume $\mathcal{A}/2\pi = 200$ MHz in all results shown in the following slides. All other parameters are the same as in page 2.



Three-level model of the qubit



t(ns)





















